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S/057/61/031/002/002/015
B124/B204

The excitation of electromagnetic ...

$$\begin{aligned} \epsilon_{12} &= - \sum_{\alpha} \frac{2\pi Q_{\alpha}^2}{\omega \omega_{\alpha}} \sum_{\beta} \int v_{\perp} dv_{\perp} v_{\parallel} dv_{\parallel} \frac{R_{\alpha} J_{\beta}}{s(s+b)}, \\ \epsilon_{22} &= \sum_{\alpha} \frac{2\pi Q_{\alpha}^2}{\omega \omega_{\alpha}} \sum_{\beta} \int v_{\perp} dv_{\perp} v_{\parallel} dv_{\parallel} \frac{R_{\alpha} J_{\beta}'}{s(s+b)}, \\ \Omega_{\alpha}^2 &= \frac{4\pi e^2 n_{0\alpha}}{m_{\alpha}}, \quad \omega_{\alpha} = \frac{c k_{\parallel}}{m_{\alpha}}, \quad a = \frac{k_{\perp} v_{\perp}}{\omega_{\alpha}}, \quad b = \frac{k_{\parallel} v_{\parallel} - \omega}{\omega_{\alpha}}, \\ R_{\alpha} &= \frac{v_{\perp}}{\omega_{\alpha}} \left(-b \frac{\partial f_{\alpha}}{\partial v_{\perp}} + \text{ctg} \theta \alpha \frac{\partial f_{\alpha}}{\partial v_{\parallel}} \right), \end{aligned} \quad (2.3)$$

$J_0(a)$ is a Bessel function, $J'_0 = J'_0(a)$ is its derivative, $n_{0\alpha}$ and $f_{0\alpha}$ are the density and equilibrium function of the distribution of the particles of the kind α . The functions $f_{0\alpha}$ depend on v_{\perp} and v_{\parallel} (v_{\perp} and v_{\parallel} are the components of particle velocity which are perpendicular and/or parallel to H_0). The summation in (2.3) is carried out over all kinds of beam particles.

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in the plasma. The integration of v_{\parallel} is from $-\infty$ to $+\infty$. Eq. (2.3) (2.5)
may be obtained also when using the general equation for ϵ_{ij} (obtained by
V. D. Shafranov). The tensor ϵ_{ij} is represented in the form
 $\epsilon_{ij} = \epsilon_{ij}^{(0)} + \epsilon'_{ij}$ (2.4), where $\epsilon_{ij}^{(0)}$ denotes the tensor of the dielectric
constant of the "cold" plasma, whose components are given by (2.5):

$$\epsilon'_{11} = \epsilon'_{22} = 1 - \sum_j \frac{\omega_{pj}^2}{\omega^2 - \omega_{cj}^2}, \quad \epsilon'_{33} = 1 - \sum_j \frac{\omega_{pj}^2}{\omega^2 - \omega_{cj}^2},$$

$$\epsilon'_{12} = -\sum_j \frac{\omega_{pj}^2 \omega_{cj}}{(\omega^2 - \omega_{cj}^2)^2}, \quad \epsilon'_{13} = \epsilon'_{23} = 0. \quad (2.5)$$

The summand ϵ'_{ij} is due to the presence of the beam. Should the inequalities
 $k_{\parallel}^2 (v_{\parallel} - v_0)^2 \ll (\omega - \omega_{H\alpha} - k_{\parallel} v_0)^2$ and $k_{\perp} v_0 \ll |\omega_{H\alpha}|$ hold, if $v_0 = \bar{v}_{\parallel}$ is the drift
velocity of the beam, the thermal motion of the particles need not be taken
into account, where (2.6)

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$$\left. \begin{aligned} \epsilon_{11} = \epsilon_{22} &= - \sum_j \frac{\Omega_{pj}^2 (\omega - k_{\parallel} v_{pj})^2}{\omega^2 [(\omega - k_{\parallel} v_{pj})^2 - \omega_{pj}^2]} \\ \epsilon_{12} &= - \sum_j \frac{\Omega_{pj}^2}{\omega^2} \left[\frac{\omega^2}{(\omega - k_{\parallel} v_{pj})^2} + \frac{(k_{\perp} v_{pj})^2}{(\omega - k_{\parallel} v_{pj})^2 - \omega_{pj}^2} \right] \\ \epsilon_{21} &= - \sum_j \frac{\Omega_{pj}^2 (\omega - k_{\parallel} v_{pj}) \omega_{pj}}{\omega^2 [(\omega - k_{\parallel} v_{pj})^2 - \omega_{pj}^2]} \\ \epsilon_{22} &= - \sum_j \frac{\Omega_{pj}^2 (\omega - k_{\parallel} v_{pj}) k_{\perp} v_{pj}}{\omega^2 [(\omega - k_{\parallel} v_{pj})^2 - \omega_{pj}^2]} \\ \epsilon_{33} &= \sum_j \frac{\Omega_{pj}^2 \omega_{pj} k_{\perp} v_{pj}}{\omega^2 [(\omega - k_{\parallel} v_{pj})^2 - \omega_{pj}^2]} \end{aligned} \right\} \quad (2,6)$$

and the cross bar denote that the respective quantity refers to the beam.
For studying the effect of thermal leakage of the beam particles by ex-
citement of the waves of the plasma, the equilibrium function of the beam

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particle distribution is chosen in the form

$$f_{0\alpha} = \frac{1}{(2\pi)^{3/2} v_{T\alpha}^3} \exp \left[-\frac{v_1^2 + (v_1 - v_0)^2}{2v_{T\alpha}^2} \right], \quad v_{T\alpha}^2 = \frac{T_\alpha}{m_\alpha} \quad (2.7)$$

where the tensor ϵ'_{ij} assumes the form

$$\epsilon'_{11} = \sum_{\alpha} \frac{q_\alpha^2}{k^2} x_\alpha e^{-x_\alpha^2} \left(\frac{x_\alpha^2}{k^2} + 2x_\alpha \right) / \sqrt{\pi} w(x_\alpha),$$

$$\epsilon'_{22} = \sum_{\alpha} \frac{q_\alpha^2}{k^2} x_\alpha e^{-x_\alpha^2} \left[\left(\frac{x_\alpha^2}{k^2} + 2x_\alpha \right) / k - 2x_\alpha \right] / \sqrt{\pi} w(x_\alpha),$$

$$\epsilon'_{33} = \sum_{\alpha} \frac{q_\alpha^2}{k^2} \left[x_\alpha^2 + \sum_{\alpha} x_\alpha e^{-x_\alpha^2} / w(x_\alpha) \right],$$

$$\epsilon'_{12} = \sum_{\alpha} \frac{q_\alpha^2}{k^2} x_\alpha e^{-x_\alpha^2} (x_\alpha - k) \sqrt{\pi} w(x_\alpha),$$

$$\epsilon'_{13} = \sum_{\alpha} \frac{q_\alpha^2}{k^2} x_\alpha e^{-x_\alpha^2} / \sqrt{\pi} w(x_\alpha),$$

$$\epsilon'_{23} = \sum_{\alpha} \frac{q_\alpha^2}{k^2} x_\alpha e^{-x_\alpha^2} \sqrt{\pi} (x_\alpha - k) \sqrt{2\pi} w(x_\alpha),$$

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where $I_s = I_s(\mu)$ is a modified Bessel function, $I'_s = dI_s/d\mu$,

$$\sqrt{\mu} = k_{\perp} v_{h,n} / \omega_{H\alpha}, \quad z_s = (\omega - s\omega_{H\alpha} - k_{\parallel} v_0) / \sqrt{2} k_{\parallel} v_{h\alpha}, \quad \nu_s = (\omega - s\omega_{H\alpha}) / \sqrt{2} k_{\parallel} v_{h\alpha},$$

and $w(z_s)$ is the probability integral. In the further course of the work, equations are derived for the electron longitudinal oscillations, the quasi-longitudinal propagation as well as the ion-cyclotron and magnetohydrodynamic waves. Mention is made of A. I. Akhiezer, V. V. Zheleznyakov, Ya. B. Faynberg, and V. P. Dokuchayev. A. I. Akhiezer and V. P. Aleksin are thanked for discussion and valuable advice. There are 12 references: 10 Soviet-bloc and 2 non-Soviet-bloc.

ASSOCIATION: Fiziko-tekhnicheskii institut AN USSR, Khar'kov
(Institute of Physics and Technology AS UkrSSR, Khar'kov)

SUBMITTED: May 9, 1960

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26.2330
24.2120(1482, 1502, 1160)
AUTHORS: Kitsenko, A. B. and Stepanov, K. N.

TITLE: Cyclotron instability in plasma

PERIODICAL: Zhurnal tekhnicheskoy fiziki, v. 31, no. 2, 1961, 176-179

TEXT: It is known that the anisotropy of the equilibrium distribution of plasma electrons and ions with respect to velocity may lead to instabilities. The present paper deals with the instability of an unbounded plasma in a homogeneous magnetic field H_0 if ions with an equilibrium distribution function of the form (1)

$$f_i = \frac{n_i}{\pi^{3/2}} \exp \left(-\frac{v_{\perp}^2}{v_{\perp 0}^2} - \frac{v_{\parallel}^2}{v_{\parallel 0}^2} \right)$$

occur in the plasma, where n_i the density of the ions, v_{\perp} and v_{\parallel} the perpendicular and parallel components, respectively, of the ion velocity. The ions with the distribution function (1) have the same Larmor radius $r_0 = v_{\perp 0} / \omega_H$, where $\omega_H = eH_0 / Mc$ and M the mass of the ion. It is demonstrated

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that disturbances of a wavelength of the order of r_0 and frequencies which are equal to or a multiple of the gyrofrequency ω_H are unstable (cyclotron instability). The order of magnitude of the growth increment is equal for the first harmonics. It may be expected that the results obtained for an unbounded plasma will be qualitatively also correct for adiabatic traps. First, the case is studied where the thermal scattering of ion velocities parallel to H_0 can be neglected. Hence, it may be assumed that $f_0(v_{||}) = \delta(v_{||})$. The tensor of the dielectric constant in this case has

$$\text{form (2), } \epsilon_{11} = -\frac{Q^2}{\omega^2} \left[1 + \sum_{n=-\infty}^{\infty} \left(\frac{2\pi n^2 J_n^2}{(\omega - n\omega_H)^2} + \frac{c_1^2 \omega_H^2 J_n^2}{(\omega - n\omega_H)^4} \right) \right],$$

$$\epsilon_{22} = -\frac{Q^2}{\omega^2} \left[1 + \sum_{n=-\infty}^{\infty} \left(\frac{\omega_H n (c_1^2 J_n^2)'}{(\omega - n\omega_H)^2} + \frac{c_1^2 \omega_H^2 J_n^2}{(\omega - n\omega_H)^4} \right) \right], \quad \left. \begin{matrix} (2) \\ a) \end{matrix} \right\}$$

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$$\begin{aligned} \epsilon_n &= 1 - \frac{\Omega^2}{\omega^2}, \\ \epsilon_{11} &= -\frac{\Omega^2}{\omega^2} \sum_{n=-\infty}^{\infty} \left(\frac{u_{en}^2 (a J_n' / J_n)}{(\omega - nu_e)^2} + \frac{u_{en}^2 \omega_{en} J_n' / J_n}{(\omega - nu_e)^3} \right) - \frac{\Omega^2}{\omega^2}, \\ \epsilon_{12} &= -\frac{\Omega^2}{\omega^2} \sum_{n=-\infty}^{\infty} \frac{u_{en} \omega_{en} J_n' / J_n}{\omega - nu_e}, \\ \epsilon_{22} &= -\frac{\Omega^2}{\omega^2} \sum_{n=-\infty}^{\infty} \frac{u_{en} \omega_{en} J_n' / J_n}{\omega - nu_e}, \end{aligned} \quad (2)$$

where $\Omega = (4\pi e^2 n_0 / M)^{1/2}$ the Langmuir ion frequency, θ the angle between the wave vector \vec{k} and \vec{H}_0 ; $k_{||} = k \cos \theta$; $k_{\perp} = k \sin \theta$; $a = k_{\perp} r_0$; $I_n = I_n(a)$ a Bessel function; the prime denotes the differentiation with respect to a ; $\omega = \omega_r + i\gamma$ is the complex frequency. The axis Oz is parallel to \vec{H}_0 , the axis Ox lies in the plane of the vectors \vec{k} and \vec{H}_0 . If the electrons have

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Maxwellian velocity distribution, (3) holds

$$\left. \begin{aligned} \epsilon_{ij} &= \frac{\omega_p^2}{k_{\parallel}^2 v_{Te}^2} [1 + i \sqrt{\pi} w(x)], \quad v_{Te}^2 = \frac{T_e}{m}, \\ w(x) &= e^{-x^2} \left(\frac{-k_{\parallel}}{|k_{\parallel}|} + \frac{2i}{\sqrt{\pi}} \int_0^x e^{t^2} dt \right), \quad x = \frac{\omega - \omega_{ce}}{\sqrt{2} k_{\parallel} v_{Te}}, \end{aligned} \right\} \quad (3)$$

where Ω_e the Langmuir electron frequency. The solutions are then obtained for the dispersion relations for electromagnetic waves near $\omega_s = s \Omega_e$ in the plasma (Ref. 1):

$$\omega' = \omega_s + \xi, \quad |\xi| \ll \omega_s, \quad s = 0, \pm 1, \pm 2, \quad (4)$$

If the quantity $a = k_{\perp} r_D$ is not too low and the number of the harmonics s not too high, the last summands for ϵ_{ij} in equation (2) $\sim 1/(\omega' - s \Omega_e)^2$ are the highest. Taking account of this fact

$$\epsilon_{ij} = - \frac{\omega_p^2}{k_{\perp}^2 v_{Te}^2} (e^{i\theta} J_0^2 + \cos^2 \theta e^{i\theta} J_1^2) + \frac{\omega_p^2 k_{\perp}^2 J_0^2}{\omega_s^2 [1 + i \sqrt{\pi} w(x)]}, \quad (5)$$

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holds, where $z_0 = \omega / \sqrt{2} k_{||} v_{te}$. If $a \sim 1$, $|z_0| \sim v_0 / v_{te}$. If it appears from (5) that $\gamma = \text{Im} z > 0$, the plasma is unstable. This is always the case when the first summand in (5) is greater than the second one, and also for $|z_0| > 1$, if (6).

$$\gamma = -\frac{1}{2} \frac{\partial^2 f_0}{\partial v^2} \left(v^2 f_0' + \cos^2 \theta v f_0' \right) + \frac{\partial^2 f_0}{\partial v^2}$$

Finally the case is dealt with where the thermal motion of the ions is essential. For reasons of simplicity, function $f_0(v_{||})$ is assumed to be of Maxwellian type, i.e., (7)

$$f_0(v_{||}) = \frac{1}{\sqrt{2\pi} v_{te}} \exp \left(-\frac{v_{||}^2}{2v_{te}^2} \right), \quad v_{te}^2 = \frac{T_e}{M}$$

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In this case, the tensor ϵ_{ij} is determined by equations (8),

$$\begin{aligned} \epsilon_{11} &= -\frac{\omega_p^2}{\omega^2} \left[1 - \frac{\omega_p^2}{\omega^2} \sum_k \left(\sqrt{\pi} \omega_k(x_k) \left(\frac{\omega_k}{\omega} n^2 J_0^2 + 2n^2 J_0 J_1 \frac{\omega_k}{\omega} \right) \right) \right], \\ \epsilon_{22} &= -\frac{\omega_p^2}{\omega^2} \left[1 - \frac{\omega_p^2}{\omega^2} \sum_k \left(\sqrt{\pi} \omega_k(x_k) \left(\frac{\omega_k}{\omega} J_0^2 + \frac{\omega_k}{\omega^2} (n^2 J_1^2) \right) \right) \right], \\ \epsilon_{33} &= 1 + \frac{\omega_p^2}{\omega^2} \left[2x_0^2 + \frac{\omega_p^2}{\omega^2} + x_0^4 + \right. \\ &\quad \left. + \sum_k 2\sqrt{\pi} \omega_k(x_k) \left(x_0 J_0 + \frac{2x_0 \omega_k}{\omega^2} n J_1 \right) \right], \\ \epsilon_{12} &= -\frac{\omega_p^2}{\omega^2} \sum_k \sqrt{\pi} \omega_k(x_k) \left[\frac{\omega_k}{\omega^2} (n J_1) + \frac{\omega_k}{\omega} n \omega J_0 \right] - \frac{\omega_p^2}{\omega^2}, \\ \epsilon_{13} &= \frac{\omega_p^2}{\omega^2} \left[1 - \frac{\omega_p^2}{\omega^2} + \sum_k \left(\sqrt{\pi} \omega_k(x_k) \left(\frac{2n^2 J_0 J_1}{\omega} + \frac{\omega_k}{\omega^2} J_0^2 \right) \right) \right], \\ \epsilon_{23} &= \frac{\omega_p^2}{\omega^2} \sum_k \left(\sqrt{\pi} \omega_k(x_k) \left[\frac{\omega_k}{\omega} (n J_1) + \frac{\omega_k}{\omega^2} J_0 \right] \right). \end{aligned} \quad (8)$$

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The authors studied the occurrence of the instability for frequencies near ω_s . As may be seen from (5), the growth increment is small compared with the frequency if $v_0 \leq v_A$. For high-density plasmas, where $\Omega \gg \omega_H$, stronger instabilities occur in the case of disturbances with a wavelength of the order of r_0 , for which growth frequency and increment are in the order of ω_H . In this case ω and γ can be determined only numerically. There is 1 Soviet-bloc reference.

ASSOCIATION: Fiziko-tekhnicheskiy institut AN USSR, Khar'kov (Institute of Physics and Technology of the AS UkrSSR, Khar'kov)

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9.2572 (a/w 1532)

10.2000

24.6716

AUTHOR:

Kitsenko, A. B.

TITLE:

Character of the instability in the interaction of a beam with plasma in an outer magnetic field

PERIODICAL: Zhurnal tekhnicheskoy fiziki, v. 31, no. 10, 1961, 1270-1271

TEXT: A convective instability can only be eliminated if its character is known. This article deals with the excitation of h-f longitudinal electron oscillations during the passage of a charged particle beam through plasma in an outer magnetic field. The resulting instability is convective in the absence of a magnetic field (Ref. 3, see below). A similar instability occurs in an outer magnetic field due to cyclotron excitation on the basis of the anomalous Doppler effect. This instability is investigated in the following. Integrals of the form $\int a(k) \cdot 1k^2 - 10^2 dk$ describe the electric and the magnetic field strength and can be used to distinguish between the absolute and the convective instability. If this integral tends to zero for any constant \bar{r} with $t \rightarrow \infty$, the instability will be convective, and if it tends to infinity, it will be absolute. The integral

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Character of the instability in the

may also be written as $\int_a[k(\omega)] \frac{d\omega}{dk} d\omega$ where C is the mapping of

the real axis of the k -plane on the ω -plane. The values of ω for

which $d\omega/dk = 0$, are found from the dispersion equation to be

$$\omega_0^2 = \omega_H^2(1 + \eta) \pm \sqrt{(1 + \eta)^2 - 4\cos^2\theta\eta}/2\eta$$

$$\eta = \frac{\omega_H^2 - \Omega^2}{\Omega^2}, \text{ where } \omega_H \text{ is the electron cyclotron frequency; } \Omega \text{ and } \Omega_0$$

are the Langmuir frequencies of plasma and beam; θ is the angle between H_0 and K . The dispersion equation shows that C cuts the imaginary axis

only in $\omega = 0$. The contour C can be deformed such that it lies in the

lower semiplane ω , i.e., that the integral for $t \rightarrow \infty$ tends to zero.

K. N. Stepanov, R. V. Polovin, and V. D. Shapiro are thanked for discussions. There are 5 references: 3 Soviet and 2 non-Soviet. Ref. 3:

Ya. B. Faynberg et al., ZhTF, 31, 633, 1961. The two references to English-language publications read as follows: Ref. 1: P. A. Sturrock, Phys. Rev., 112, 1488, 1958; Ref. 2: P. A. Sturrock, Phys. Rev., 117.

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Character of the instability in the ...
1426, 1960.

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technical Institute, AN UkrSSR, Khar'kov)

SUBMITTED: February 28, 1961

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B125/B202

24.6714

AUTHORS: Akhiezer, A. I., Kitsenko, A. B., Stepanov, K. N.

TITLE: Interaction between charged particle currents and low-frequency plasma oscillations

PERIODICAL: Zhurnal eksperimental'noy i teoreticheskoy fiziki, v. 40, no. 6, 1961, 1866-1870

TEXT: The authors deal with the interaction between a compensated beam of charged particles and the low-frequency oscillations of a plasma (mainly with the magneto-acoustic waves and the Alfvén waves) in a constant field in parallel direction to the beam and in the absence of collisions. If the plasma is rarefied to such an extent that the frequency ω of the oscillations is much higher than the frequency $1/\tau$ of the collisions, the plasma oscillations must be described on the basis of the kinetic equation. With $\omega\tau \ll 1$ the plasma can be described hydrodynamically. The authors studied the case $\omega\tau \gg 1$. The general dispersion equation for plasma oscillations in an external magnetic field with random distribution function of the particles with respect to the velocities

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reads as follows: $An^4 + Bn^2 + C = 0$ (1) where $n = k\omega/\omega$. The wave vector \vec{k} and the quantities A, B, and C are determined by the components of the tensor of the dielectric constant ϵ_{ij} . Furthermore, it is assumed that $\omega \ll \omega_{H1}$, $kv_0 \ll \omega_{H1}$ where ω_{H1} is the gyrofrequency of the ions, v_1

$-(T_1/M)^{1/2}$ the mean thermal velocity of the ions (T_1 denotes the temperature and M the mass of the ions) and v_0 the velocity of the beam. Under these conditions (1) falls into the equations $(k\omega/\omega)^2 \cos^2 \theta - \epsilon_{11} = 0$ (3) and $(k\omega/\omega)^2 - \epsilon_{22} - \epsilon_{23}^2/\epsilon_{33} = 0$ (4) describing the Alfvén wave and the sound wave, respectively. If the velocity distribution of the particles

in the beam has the form
$$f_{0,i} = n_0 \left(\frac{m_{e,i}}{2\pi T_{e,i}} \right)^{3/2} \exp \left\{ -\frac{m_{e,i}(v-v_0)^2}{2T_{e,i}} \right\} \quad (5)$$

(n_0 density of the particles in the beam, T_0 , T_1 temperatures of the electrons and ions of the beam, $m_e = m$, $m_1 = M$)

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$$\epsilon_{11} = 1 + \sum_a \frac{\Omega_a^2 (\omega - k_1 v_{0a})^2}{\omega_{H0}^2 \omega^2}, \quad \epsilon_{21} = \epsilon_{11} + \sum_a \frac{\Omega_a^2 H_0^2}{\omega_{H0}^2 \omega^2} 2i \sqrt{\pi} \sin^2 \theta z_a w(z_a),$$

$$\epsilon_{33} = 1 + \sum_a \frac{\Omega_a^2}{k_1^2 v_a^2} (1 + i \sqrt{\pi} z_a w(z_a)),$$

$$\epsilon_{23} = - \sum_a \frac{\Omega_a^2}{\omega \omega_{H0}} \sqrt{\pi} i g(z_a) w(z_a), \quad (6)$$

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$$w(z_a) = e^{-z_a^2} \left(\pm 1 + \frac{2i}{\sqrt{\pi}} \int_0^{z_a} e^{t^2} dt \right), \quad z_a = \frac{\omega - k_1 v_{0a}}{\sqrt{2} k_1 v_a},$$

$$\Omega_a^2 = 4\pi e^2 n_{0a} / m_a, \quad v_a^2 = T_a / m_a, \quad \omega_{H0} = e H_0 / m_e c, \quad k_1 = k \cos \theta$$

holds with Maxwellian equilibrium velocity distribution of the electrons and ions of the plasma. The upper and lower signs in $w(z_a)$ hold with $k_{||} > 0$ and $k_{||} < 0$, respectively; the index a denotes the types of all particles of the plasma and the beam. With the aid of (6) expression

$$\omega = k_1 \frac{v_0 \Omega_i^2 \pm ((\Omega_i^2 + \Omega_e^2) e^{-\frac{1}{2} \frac{\omega_{H0}^2}{\omega^2}} - \Omega_i^2 \Omega_e^2 v_e^2)^{1/2}}{\Omega_i^2 + \Omega_e^2}, \quad (7)$$

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for the frequency of the Alfvén wave modified by the existence of the beam is obtained from (3). Due to the excitation of the Alfvén waves the state of the system plasma - beam is unstable if condition

$$\omega^2 > V_A^2 + V^2, \quad (8)$$

$$V_A = H_0 / \sqrt{4\pi n_0 M}, \quad V = H_0 / \sqrt{4\pi n_0' M}.$$

is fulfilled. With sufficiently low densities of the beam the excitation of the Alfvén waves is impossible as long as the coupling between the Alfvén waves and the magneto-acoustic waves is neglected. When considering this coupling an instability is observed also with those densities at which (8) is not valid. In the following study of the beam by means of the magneto-acoustic waves the density of the beam is assumed to be small as compared to the plasma density. With $kv_1 \ll \omega \ll kv_0$ the solution of the dispersion equation (4) has the form

$$\omega_{\pm} = \pm V_A, \quad V_A^2 = \frac{1}{2}(V_0^2 + s^2 \pm [(V_0^2 + s^2)^2 - 4V_0^2 s^2 \cos^2 \theta]^{1/2}). \quad (9)$$

with lacking beam, where $s = (T_0/M)^{1/2}$. With $V_A \sim s$ the quantities V_{\pm} are also of the order of magnitude s . In this case the condition $\omega \gg kv_1$

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holds only for a strongly nonisothermal plasma ($T_e \gg T_i$) while the magneto-acoustic waves with $T_e \lesssim T_i$ are strongly attenuated. Besides, also $|\omega - k_{\parallel} v_0| \gg kv_0$ holds. The dispersion equation (4) then reads as follows: $\omega = k_{\parallel} v_0 + \xi$ (10) with $|\xi| \ll |k_{\parallel} v_0|$. Under the conditions studied, the state of the system plasma-beam is unstable due to the excitation of the magneto-acoustic waves. If v_0 does not lie in the interval $s < v_0 < v_A$ this instability occurs even with degeneration of η ($|\eta| \ll 1$). With $v_0 \cos \theta \rightarrow v_{\perp}$, (13) holds. With the maximum (resonance) interaction (with $v_{\perp} = v_0 \cos \theta$) the increment of increase of the oscillations is not proportional to $(n'_0/n_0)^{1/2}$ but to $(n'_0/n_0)^{1/3}$. The solutions (10) to (13) of the dispersion equation (4) hold for a strongly nonisothermal plasma. With $T_e \lesssim T_i$,

$$s = \left(\frac{M}{m}\right)^{1/2} s_0 = \pm \frac{\Omega_e (n^2 - s_{\perp}^{(0)})^{1/2}}{[s_{\perp}^{(0)} (n^2 - s_{\perp}^{(0)}) - s_{\perp}^{(0)2}]^{1/2}} \quad (14)$$

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Interaction between charged ...

is obtained with neglect of the thermal motion of the electrons of the beam. ϵ_{1j}^0 is the component of the tensor of the dielectric constant of the plasma without beam with $\omega = k_{\parallel} v_0$. An instability also occurs with $T_e \lesssim T_i$ if the plasma oscillations are weakly attenuated. With $k v_i \ll |\omega - k_{\parallel} v_0| \ll k v_0$ the thermal motion of the electrons in the beam has to be taken into account. (4) then has the solution $\omega = k_{\parallel} v_0 + \epsilon_0$, $|\epsilon_0| \ll k_{\parallel} v_0$ (15) where ϵ_0 is obtained from (14). If $k_{\parallel} v_0$ lies near the eigenfrequency $k v_{\pm}$ of the magneto-acoustic wave in the nonisothermal plasma $\omega = k v_{\pm} + \epsilon_0$, $|k(v_0 \cos \theta - v_{\pm})| \ll |\epsilon_0|$ (16) holds where ϵ_0 is to be determined from (13). These formulas hold for sufficiently low temperatures of the beam $|\omega - k_{\parallel} v_0| \gg k v_i$. With sufficiently small n_0'/n , $\omega = \omega_{\pm} + i\gamma_{\pm}$ holds with

$$\gamma_{\pm} = -\frac{\sqrt{n_0'} \sin^2 \theta (\epsilon_+ + \epsilon_0 + \epsilon_i)}{4\epsilon_{\pm} |\cos^2 \theta - \epsilon_{\pm}| (\epsilon_{\pm} - \epsilon_0)}, \quad \epsilon_{\pm} = \left(\frac{v_{\pm}}{v_0}\right)^2.$$

Card 6/7

25206

S/056/61/040/006/027/031

B125/B202

Interaction between charged ...

A beam with low density and high energy spread of the electrons and ions generally does not cause a magneto-acoustic wave in the plasma. There are 9 references: 7 Soviet-bloc and 2 non-Soviet-bloc. The two most recent references to English-language publications read as follows:
D. Bohm, E. Gross. Phys.Rev., 75, 1851, 1864, 1949. I.B. Bernstein.
R.M. Kulsrud. Phys.Fl., 3, 937, 1960.

ASSOCIATION: Fiziko-tehnicheskii institut Akademii nauk Ukrainasoy SSR
(Institute of Physics and Technology of the Academy of
Sciences of the Ukrainskaya SSR)

SUBMITTED: January 27, 1961

Card 7/7

35355

3/057/62/032/003/006/019
B116/B102

24.6714
AUTHORS:

Kitsenko, A. B., and Stepanov, K. N.

TITLE:

Excitation of magnetosonic waves in dilute plasma by a charged particle flow

PERIODICAL: Zhurnal tekhnicheskoy fiziki, v. 32, no. 3, 1962, 303 - 307

TEXT: The excitation of "fast" and "slow" magnetosonic waves in dilute plasma by a charged particle flow of arbitrary density was studied. The following assumptions were made: The Alfvén velocity is much greater than the sonic velocity, plasma and particle flow are dilute so that "near" collisions may be neglected, space charge and electric current of the particle flow are compensated in the state of equilibrium. The dispersion equation for magnetosonic waves

$\frac{k^2 c^2}{\omega^2} - \epsilon_{22} - \frac{\epsilon_{23}^2}{\epsilon_{33}} = 0$ (1) is written down, where ϵ_{ij} is the tensor of the dielectric constant (axis 3 runs parallel to the external magnetic field H_0 , the wave vector k is in the 1-3 plane). (1) holds if $\omega \ll \omega_{H1}$.

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S/057/62/032/003/006/019
B116/B102

Excitation of magnetosonic...

$kv_1 \ll \omega_{H1}$, $kv_0 \ll \omega_{H1}$, where ω_{H1} and v_1 = hydrogen frequency and mean thermal velocity of ions respectively, v_0 = beam velocity ($\vec{v}_0 \parallel \vec{H}_0$). If beam and plasma particles show Maxwellian velocity distribution, the components of ϵ_{ij} are:

$$\begin{aligned}\epsilon_{11} &= 1 + \sum_j \frac{\Omega_j^2 (v - k_1 v_j)^2}{\omega_{Hj}^2 \omega_j^2}, \\ \epsilon_{22} &= \epsilon_{11} + \sum_j \frac{\Omega_j^2 k^2 v_j^2}{\omega_{Hj}^2 \omega_j^2} 2/\sqrt{\pi} \sin^2 \theta z_j w(z_j), \\ \epsilon_{33} &= 1 + \sum_j \frac{\Omega_j^2}{k_1^2 \omega_j^2} (1 + i/\sqrt{\pi} z_j w(z_j)), \\ \epsilon_{23} &= - \sum_j \frac{\Omega_j^2}{\omega_{Hj} \omega_j} i g \theta \sqrt{\pi} z_j w(z_j), \\ w(z_j) &= e^{-z_j^2} \left(\pm 1 + \frac{2i}{\sqrt{\pi}} \int_0^{z_j} e^{t^2} dt \right), \quad z_j = \frac{v - k_1 v_j}{\sqrt{2} k_1 v_j}, \\ \Omega_j^2 &= \frac{4\pi e^2 n_{0j}}{m_j}, \quad v_j^2 = \frac{T_j}{m_j}, \quad \omega_{Hj} = \frac{e H_0}{m_j c}, \quad k_1 = k \cos \theta.\end{aligned}\tag{2}$$

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3/057/62/032/003/006/019
B116/B102

Excitation of magnetosonic...

θ is the angle between \vec{k} and \vec{H}_0 . First, the oscillations corresponding to the "fast" magnetosonic wave are investigated. It is assumed that $k \gg v_1, v_1', v_0, m/M, v_0', m/M$, the indices e and i denote electrons and ions respectively, the primes refer to beam particles. The solution of (1) is $\omega = \omega_{\pm}$, where

$$\omega_{\pm} = k \frac{a_e a_i \cos \theta \pm \sqrt{a_e^2 (V_A^2 + V_A'^2 - a_0^2 \cos^2 \theta)}}{a_e + a_i} \quad (3)$$

$$V_A = \frac{H_0}{\sqrt{4\pi n_0 M}}, \quad V_A' = \frac{H_0}{\sqrt{4\pi n_0' M}}.$$

From (3) follows the instability condition $v_0^2 > V_A^2 + V_A'^2$. Consideration of the kinetic effects shows that instability may also occur at $v_0^2 < V_A^2 + V_A'^2$. In this case, $\omega = \omega_{\pm} + i\gamma_{\pm}$, $|\gamma_{\pm}| \ll |\omega_{\pm}|$. On the above assumptions, the dispersion equation for "slow" magnetosonic waves reads $\epsilon_{33}(\omega, k) = 0$ (6). For the case $|\omega - k_{||} v_0| \gg kv_0'$, (6) gives

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Excitation of magnetosonic...

S/057/62/032/003/006/019
B116/B102

$$\frac{\omega_i^2}{\omega^2} + \frac{\omega_i^2}{(\omega - k_1 v_0)^2} = \frac{\omega_i^2}{k_1^2 s^2} (1 + i\sqrt{\pi} z), \quad (7)$$

where $s = T_e/M$ is the sonic velocity, $z_0 = \frac{\omega}{\sqrt{2} |k_1| v_0}$ ($|z_0| \ll 1$). If $z_0 = 0$,
the instability condition $v_0^2 < \frac{s^2}{n_0} \left[n_0^2 + \left(n_0 \frac{M}{\pi} \right)^2 \right]$. (8)

is obtained. Explicit formulae for the increments γ can be obtained from
(7) for some limiting cases only. For the case $kv_1^2 \ll |\omega - k_{||} v_0| \ll kv_0^2$,
(6) gives

$$\frac{\omega_i^2}{\omega^2} + \frac{\omega_i^2}{(\omega - k_1 v_0)^2} = \frac{\omega_i^2}{k_1^2 s^2} (1 + i\sqrt{\pi} z) + \frac{\omega_i^2}{k_1^2 s^2} (1 + i\sqrt{\pi} z). \quad (12)$$

Neglecting z_1 and z_2 , the instability condition

$$v_0^2 < \frac{s^2 n_0^2 (2(n_0^{1/3} + n_0^{-1/3}))^2}{n_0 s^2 + n_0 s^2} \quad (13).$$

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Excitation of magnetosonic...

S/057/62/032/003/006/019
B116/B102

is obtained. The solution of (12) can be found explicitly in some special cases. For $|\omega - k v_0| \ll k v_1$, (7) gives $\omega = k u + i\gamma$, where

$$\frac{1}{\omega} = \frac{1}{k} + \frac{\gamma}{k^2} \left(\frac{1}{u} + \frac{1}{v_1} \right),$$

$$\gamma = -\sqrt{\frac{\pi}{2}} v_1 \left\{ \sqrt{\frac{\pi}{2}} \frac{\gamma}{k} + (u - v_1) \frac{\gamma}{k^2} \left(\sqrt{\frac{\pi}{2}} \frac{1}{u} + \frac{1}{v_1} \right) \right\} \quad (18)$$

Then the instability condition is $v_0 - u > \frac{\gamma v_1^2}{k^2 \left(\sqrt{\frac{\pi}{2}} \frac{1}{u} + \frac{1}{v_1} \right)}$. (19)

A. I. Akhiezer is thanked for advice. There are 7 references: 6 Soviet and 1 non-Soviet. The reference to the English-language publication reads as follows: I. B. Bernstein, R. M. Kulsrud. Phys. Fl., 3, 937, 1960.

SUBMITTED: January 31, 1961 (initially)
May 3, 1961 (after revision)

Card 5/5

KITSENKO, A.B.

Interaction of charges with an electron plasma in a magnetic field.
Dokl.AN SSSR 145 no.2:305-308 J1 '62. (MIRA 15:7)

1. Fiziko-tekhnicheskiy institut AN USSR. Predstavleno akademikom
M.A.Leontovichem.
(Plasma (ionized gases)) (Magnetic fields)

ACCESSION NR: AT4036038

S/2781/63/000/003/0003/0017

AUTHORS: Kitsenko, A. B.; Stepanov, K. M.

TITLE: Investigation of plasma oscillations in the quasihydrodynamic approximation

SOURCE: Konferentsiya po fizike plazmy* i problemam upravlyayemogo termoyadernogo sinteza. 3d, Kharkov, 1962. Fizika plazmy* i problemy* upravlyayemogo termoyadernogo sinteza (Plasma physics and problems of controlled thermonuclear synthesis); doklady* konferentsii, no. 3. Kiev, Izd-vo AN UkrSSR, 1963, 3-17

TOPIC TAGS: plasma oscillation, magnetohydrodynamics, ionized plasma, plasma ion oscillation, plasma wave reflection, terrestrial magnetism, solar corpuscular radiation, ionosphere

ABSTRACT: The longitudinal oscillations of a plasma with nonisotropic distribution function, and the excitation of magnetohydrody-

Cord. 1/3

ACCESSION NR: AT4036038

namic waves in interpenetrating plasma streams with nonisotropic particle velocity distribution functions, are considered on the basis of the quasihydrodynamic equations. It is confirmed on the basis of these equations that one-dimensional Langmuir oscillations of an electron gas can be regarded as adiabatic with an adiabatic exponent $\gamma = 3$. The effects of the ions on both high-frequency and low-frequency oscillations are considered. It is shown that anomalous dispersion occurs in both frequencies. In a plasma consisting of two species of ions, plasma resonance in the region of both high and low frequencies has the same distinguishing features, namely the limitation of the growth of the refractive index because of the presence of thermal motion of the plasma particles, and the appearance of a new branch of oscillations (plasma waves) with anomalous dispersion properties. An experimental investigation of the propagation of radio waves through a plasma at resonance makes it possible to determine several important characteristics of the plasma. The excitation of magnetohydrodynamic waves when streams

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ACCESSION NR: AT4036038

of charged particles pass through a plasma is of particular interest in connection with the interaction between solar corpuscular streams in the upper ionosphere, for the plasma instability arising in the case of such a stream can lead to variations of the earth's magnetic field and to the appearance of ultralow frequency radiation. "In conclusion the authors are deeply grateful to A. I. Akhiezer for interest in the work." Orig. art. has: 32 formulas.

ASSOCIATION: None

SUBMITTED: 00

DATE ACQ: 21May64

ENCL: 00

SUB CODE: ME, AA

NR REF SOV: 016

OTHER: 004

Card 3/3

ANGELEYKO, V.V. [Anheleiko, V.V.], KITSENKO, A.B. [Kitsenko, O.B.]

Excitation of sound waves in a weakly ionized plasma. Ukr. fiz.
zhur. 10 no.1:16-20 Ja '65. (MIRA 18:4)

1. Fiziko-tekhnicheskly institut AN UkrSSR, Khar'kov.

I. 40942-65 KPP(c)/EWT(1)/EHC(t) PI-L IJP(c) 00/WW

8/0051/65/035/003/0470/0474

ACCESSION NR: APS007293

AUTHOR: Angeleyko, V.V.; Kitsenko, A.B.

TITLE: On the cyclotron excitation of magneto-acoustic waves by a stream of charged particles

SOURCE: Zhurnal tekhnicheskoy fiziki, v.35, no.3, 1965, 470-474

TOPIC TAGS: plasma beam interaction, plasma stability, magnetic sound wave, cyclotron resonance

ABSTRACT: The authors discuss the interaction of a neutral stream of charged particles with the low frequency oscillations (fast and slow magneto-acoustic waves) of a highly anisothermal plasma in an external magnetic field. The magnetic and kinetic pressures are assumed to be of the same order of magnitude, and the interaction is discussed in the neighborhood of the cyclotron resonance under conditions of the anomalous Doppler effect. The relevant dispersion equation is quoted from earlier work (K.M.Stepanov and A.B.Kitsenko, ZhTF 31,167,1961) and is simplified for the case in which the frequency is low compared with the ion Larmor frequency and the phase velocity is high compared with the ion thermal velocity and low com-

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I. 40942-65

ACCESSION NR: AP0007292

pared with the electron thermal velocity. Solutions are obtained and discussed separately for the case of a hot beam and that of a cold beam, and for conditions of cyclotron resonance with the ions, and with the electrons of the beam. "In conclusion, the authors express their deep gratitude to A.I. Akhiezer and E.N. Stepanov for valuable advice and discussions of the work." Orig.art.has: 20 formulas

ASSOCIATION: none

SUBMITTED: 04Jun64

ENCL: 00

SUB CODE: ME

NR REF SOV: 007

OTHER: 001

Card 2/2 P45

ANGELEYKO, V.V. KITSENKO, A.B.

Cyclotron excitation of magnetoacoustic waves by a charged
particle flux. Zhur. tekhn. fiz. 35 no.3:470-474 Apr '65.
(MIRA 18:6)

L 08805-67 EWT(1) IJP(a) AT/GD
ACC NR: AT6020443 SOURCE CODE: UR/0000/65/000/000/0131/0136

AUTHOR: Kitsenko, A. B.; Gapontsev, B. A. 51

ORG: none

TITLE: Movement of streams of charged particles through a plasma at an arbitrary angle relative to the internal magnetic field

SOURCE: AN UkrSSR. Vzaïmodeystviye puchkov zaryazhennykh chastits s plazmoy (Interaction of charged particle beams with plasma). Kiev, Naukova dumka, 1965, 131-136

TOPIC TAGS: plasma magnetic field, charged particle, electron oscillation

ABSTRACT: The excitation of high frequency electron oscillations by a relativistic beam of charged particles moving in an arbitrary direction relative to the internal magnetic field is studied. The main interaction effect is assumed to take place where the wave frequency exceeds the cyclotron frequency of the beam and the wavelength is small compared with the cyclotron radius. The analysis of the tensor giving the dielectric properties of such interacting system gives the rate of growth of ordinary and extraordinary waves. The simplifying assumption which considers only waves with phase velocity much smaller than that of light yields a simple expression for the growth coefficients. In addition, so-called "whistler waves" are considered as a special case, where it is shown that small phase velocity waves are most effectively

Card 1/2

L 08805-67

ACC NR: AT6020443

generated when the angle between the wave vector and magnetic field direction approaches 90° . The case of nonrelativistic streams and their rate of growth obtained for the "whistler" mode case is also discussed. Orig. art. has: 4 figures.

SUB CODE: 20/

SUBM DATE: 11Nov65/

ORIG REF: 003

Card 2/2 net

L 08804-67 EWT(1) IJP(a) AT/QD
 ACC NR: AT6020444 (N) SOURCE CODE: UR/0000/65/000/000/0136/0143

AUTHOR: Angeleyko, V. V.; Kiteenko, A. B. 7/

ORG: none

TITLE: Excitation of longitudinal oscillations in a magnetoactive plasma

SOURCE: AN UkrSSR. Vzaimodeystviye puchkov zaryazhennykh chastits s plazmoy (Interaction of charged particle beams with plasma). Kiev, Naukova dumka, 1965, 136-143

TOPIC TAGS: magnetoactive plasma, Cerenkov radiation, Doppler effect, plasma oscillation, plasma instability

ABSTRACT: This work investigates the excitation of longitudinal oscillations in a plasma consisting of two types of ions through which a stream of particles is moving under an arbitrary angle to the direction of the internal magnetic field. The dispersion relation for the plasma is analyzed and the relationship between its four roots is studied. In the presence of the stream, a new dispersion relation is written consisting of that for the plasma and the added term for the beam and for the interaction effects. This relationship is studied for the case of a cold stream under the conditions of Cerenkov resonance and for the case of a cold stream under conditions of anomalous Doppler effect. In addition, "hot" beam interaction with the plasma is briefly discussed. The instability criteria and growth rates for the various waves is given under speci-

Card 1/2

L 08804-67

ACC NR: AT6020444

fied relationships between the frequencies characterizing the plasma and beam systems.
Orig. art. has: 23 formulas.

SUB CODE: 20/

SUBM DATE: 11Nov65/

ORIG REF: 005/

OTH REF: 002

Card 2/2 not

KITSENKO, A.V., veterinarnyy vrach

Epizootiology of anaplasmosis in cattle. Veterinariia
41 no.11:44-45 N '64. (MIRA 18:11)

1. Khorezmskaya oblastnaya veterinarnaya laboratoriya.

KITSENKO, L.

RUDITSYN, Mikhail Nikolayevich; KITSENKO, L., red.; GLEN'KAYA, I.,
tekhn.red.

[Computation and graphic work on strength of materials] Raschetno-
graficheskie raboty po soprotivleniiu materialov. Minsk. Izd-vo
Belorusakago gos.univ. im. V.I.Lenina, 1951. 227 p. (MIRA 11:7)
(Strength of materials)

ANDERS, Aleksandr Aleksandrovich; KITSENKO, M.P., inzh., retsentsent;
SMIRNOV, B.V., inzh., red.; NIKHIMSON, V.A., red. izd-va;
SOKOLOVA, T.F., tekhn. red.

[Technology of machining] Tekhnologiya mekhanicheskoi obrabotki;
sbornik zadach. Izd. 2., perer. i dop. Moskva, Gos. nauchno-tekhn.
izd-vo mashinostroit. lit-ry, 1958. 352 p. (MIRA 12:2)
(Machine-shop practice)

EWI(1)/EPA(sp)-2/T/EEC(t)/EPA(w)-2/EWA(m)-2 Pz-6/Po-4/Pab-10/P1-4

01019-65

AT

ACCESSION NR: AP50C4319

8/0185/65/010/001/0016/0020

56
58
8

AUTHOR: Anhelavko, V. V. (Anhelavko, V. V.), Kitzenko, O. (Kitzenko, A. B.)

I 31049-65

ACCESSION NR: AP5004319

authors thank A. I. Akhizer for suggesting the topic and for interest in the work, and also K. M. Stepanov and I. A. Akhizer for valuable advice." Orig. art. has: 12 formulas.

ASSOCIATION: Fiziko-tekhnicheskyy institut AN UkrSSR, Khar'kov (Physicotechnical Institute AN UkrSSR)

SUBMITTED: 02Apr64

ENCL: 00

SUB CODE: ME

NR REF SOV: 005

OTHER: 001

Card 2/2

KITSENKO, V.

In step with life. Fin. SSSR 21 no.3:61-66 Mr '60.
(MIRA 13:3)

1. Starshiy inspektor Ministerstva finansov Moldavskoy SSR.
(Kishinev--Finance)

KITSSENKO, V.P.

Morbidity among workers in the petroleum industry and methods for decreasing it. Sov. zdav. 13 no.4:21-24 J1-Ag '54. (MLRA 7:9)

**1. Is kafedry nervnykh bolezney Kubanskogo meditsinskogo instituta
(nav. prof. V.Ya. Anfinov)
(OCCUPATIONAL DISEASES, statistics,
Russia, in petroleum workers)**

KITSUNKO, V.P. (Krasnodar)

Method for treating ischias at medical centers. Kaz.-med.shur.
40 no.2:82 Mr-Apr '59. (MIRA 12:11)
(HIP JOINT--DISEASES)

KITSENKO, V.P.

Intramuscular use of novocaine under polyclinic conditions.
Vrach.delo no.7:131 JI '60.

(MIRA 13:7)

1. Glavnyy vrach polikliniki No.2 (Krasnodar).
(NOVOCAINE)

KITSENKO, V.P.

Organization of anticoagulant therapy for cardiovascular
patients under polyclinical conditions. Kas.med.shur.
no.5:83-84 8-0 '62. (MIRA 16:4)

1. Vtoraya neob'yedinennaya poliklinika Krasnodara (glavnyy
vrach - V.P.Kitsenko).

(ANTICOAGULANTS (MEDICINE))
(CARDIOVASCULAR PATIENT—CARE AND TREATMENT)

KITSENKO, V.P.; NESTEROV, V.A.

Study of industrial traumatism at the Sedin Machine-Tool
Plant. Nauch. trudy Kub. gos. med. inst. 19:56-62 '62.

(MIRA 17:8)

1. Iz kafedry organizatsii zdoravookhraneniya Kubanskogo gosudarstvennogo meditsinskogo instituta (zaveduyushchiy - dotsent V.A. Nesterov) i 2-y polikliniki g. Krasnodara (glavnyy vrach V.P. Kitzenko).

KITSENKO, V.P.

Incidence of neurological diseases in Krasnodar. Manch. trudy
Kub. gos. med. inst. 19:152-159 '62. (MIRA 17:8)

1. Iz kafedry organizatsii zdavookhraneniya (sveduyushchiy -
dokt. V.A. Nesterov) i kafedry nervnykh bolezney (sveduyushchiy -
prof. M.I. Kholodenko) Kubanskogo gosudarstvennogo meditsinskogo
instituta.

KITSSENKO, V.P.

Quality of examination and treatment of patients suffering from
neurological diseases in a polyclinic. Zdrav.Ros.Feder. 7 no.3:
21-22 Nr '63. (MIRA 1613)
(KRASNODAR--NEUROLOGY) (KRASNODAR--CLINICS)

KITSENKO, V.P.

Some problems of organizing the work of a neuropathologist in a
polyclinic. Zdrav. Ros. Feder. 7 no.9:20-22 8 '63. (MIRA 16:10)

1. Iz 2-y neob'yedinennoy polikliniki Krasnodara (glavnyy vrach
V.P. Kitzenko).

*

GOVCHAROV, S.P.; KITSENKO, V.V.; MAROULIS, A.I.; CHERNYAVSKIY, L.G.;
RZHAVSKIY, N.A.; KAMINETS, V.A., ~~redaktor~~ tekhnicheskikh nauk, redaktor; MARKUS,
M. Ye., inzhener, redaktor; MATVEYINA, Ye.N., tekhnicheskij
redaktor; SOKOLOVA, T.F., tekhnicheskij redaktor.

[Measurements of strains and stresses; handbook] Izmerenie
napriazhenii i usilii; spravochnoe posobie. Moskva, Gos. nauchno-
tekhn. izd-vo mashinostroit. lit-ry, 1955. 66 p. (MLRA 8:9)
(Strains and stresses)

KONOPLEV, Yu.V.; KITSENKO, Yu.A.; KALICHENKO, B.V.

Using the pulse neutron-neutron logging method in studying the possibility of determining the water-oil contact in horizon 4 of the Anastasiyevka-Troitskoye oil field. Geol. nefi i gaza 9 no.4:44-48 Ap '65. (MIRA 18:8)

1. Krasnodarskiy filial Vsesoyuznogo nauchno-issledovatel'skogo instituta geofizicheskikh metodov razvedki i Neftpromyslovoye upravleniye Priazovneft'.

L 33265-66 EWT(m)

ACC NR: AT6012791

(N)

SOURCE CODE: UR/3175/66/000/027/0141/0143

AUTHOR: Kitsenko, Yu. A.; Konoplev, Yu. V.

59
B+1

ORG: Krasnodar Division, VNIIGeofizika

(Krasnodarskiy filial VNIIGeofiziki)

TITLE: Introduction of the oil well impulse neutron generator IGN-1

SOURCE: USSR, Gosudarstvennyy geologicheskiy komitet. Osoboye konstruktorskoye byuro.
Geofizicheskaya apparatura, no. 27, 1966, 141-143

TOPIC TAGS: petroleum engineering,
neutron, petroleum industry equipment, temperature instrument/
IGN-1 petroleum industry equipment

ABSTRACT: This paper describes modifications of the oil well neutron impulse generator IGN-1, done locally to improve field operations under higher temperatures. To decrease heat generation, permanent magnets were substituted for the electromagnet of the ion source of the accelerating tube; oil volume was decreased by redesign, to utilize the available expansion siphone at higher temperatures; teflon insulated wires installed etc. To improve the registration effectiveness, the thermal neutron counter SNM-20 was replaced by SNM-20a, which required a higher voltage supply. Delay times, channel sensitivities and integration times were modified to increase measurement effectiveness. Tests showed stable operation in oil wells at 70°C. (Design of the IGN-1 apparatus in its entirety is not described, Abstractor). Orig.art. has 4 figures.

SUB CODE: 13, // 14
Card 1/1

SUB DATA: None/

ORIG REF: 000

VOLOVYK, T.M. [Volovyk, T.M.]; KITSENKA, L.O., [Kitseva, L.O.]

Some changes in memory block No.5 of the magnetic drum in
the electronic computer "Ural-1". Vysok L'viv. un. Ser.
matk.-mat. no.1:54-56 '65.

(MIRA 18:12)

KITSIS, A., insh.; YUZBASHOV, S., ekonomist

~~Precast~~ reinforced concrete roof panels. Stroi. mat. 4 no. 6131
Ja '58. (MIRA 11:7)

(Roofing, Concrete)

KITSIS, G.M.

Outpatient card is the medical passport of a sick person. Fel'd.
i akush. 28 no.5:50-51 №3. (MIRA 16:7)

1. Is Orgeyevskoy rayonnoy bol'nitsy, Moldavskaya SSR.
(MEDICAL RECORDS)

KITSIS, O.M.

Medical documentation at feldsher and obstetric stations. Zdravookh-
ranenie 4 no.6:42-44 N-D '61. (MIRA 15:2)

1. Is rayonnoy bol'nitsy g. Orgeyeva (glavnyy vrach Ye.M.Golovina).
(MEDICAL RECORDS)

ROZET, G.I., dotsent; ALKHAN-KEMAL, G., vrach-metodist; KITSIS, G.N.;
KRYLOV, P.M.

Letters to the editor. Zdrav.Ros.Feder. 6 no.11:35-37 N '62.
(MIRA 15:12)

1. Zaveduyushchiy kabinetom ucheta i meditsinskoy statistiki
Orgeyevskoy rayonnoy bol'nitsy Moldavskoy SSR (for Kitsis).
(PUBLIC HEALTH) (VISHNEVSKII, PETR STEPANOVICH)

COLOVINA, Ye. M.; KITSIS, G.N.

Methodological center at a district hospital. Zdravookhraneniye
6 no.5:8-9 8-0'63 (MIRA 16:12)

1. Is orgeyevskoy rayonnoy bol'nitsy (glavnyy vrach Ye.M. Golovina).

KITSIS, O.N. (g. Orgeyev)

Planning a network of public health institutions in a district.
Sovet. zdavookhr. 5:45-47 '63 (MIRA 17:2)

1. Is Orgeyevskoy rayonnoy bol'nitsy Moldavskoy SSR (glavnyy
vrach Ye.M.Colovina).

KOZLOV, Sidor Fedorovich; PAVLYUKOVICH, Aleksandr Ivanovich; KITSIS, M.S.,
red.; GRIGOR'YEVA, I.S., red. 1st-ed.; BELOGUROVA, I.A., tekhn.red.

[Hard-alloy cutting tools for the machining of light alloys] Tverdo-
splavnye' rezhushchie instrumenty dlia obrabotki legkikh splavov.
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rabotka metallov, no.25) (MIRA 16:2)
(Metal-cutting tools)

VERIGIN, N.N., prof.; Prinimali Uchastiya; KITGIS, R.A., inzh.;
ZHIGALIN, B.I., inzh.; AFINOGENOVA, N.V., inzh.;
VINOGRADOVA, G.M., red. 1st-ya; KASIMOV, D.Ya., tekhn. red.

[Methods of determining the filtration properties of rocks]
Metody opredeleniya fil'tratsionnykh svoystv gornyykh porod.
Moskva, Gos. izd-vo lit-ry po stroit., arkhit. i stroit. ma-
terialam, 1962. 177 p. (MIRA 15:4)

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snabzheniya, kanalizatsii, gidrotekhnicheskikh sooruzheniy i
inzhenernoy gidrogeologii.

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KITSIS, S., inzhener-elektromekhanik

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KITSIS, S.B.; BURDANSKIY, L.I.

Treatment of clays at the Krichev cement plant. TSement 24
no.1:24-26 Ja-Fe '58. (MIRA 11:4)

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ALESHINA, O.K., insh.; KITSIS, S.B., insh.; SHAKIMACON, N.V., insh.; KUTIN, Z.B., insh.

Using sodium fluosilicate as a mineralizer at the Krichy Cement
Factory. Nauch. soob. NIISementa no. 7:1-4 '60. (MIRA 14:5)
(Sodium fluosilicates) (Cement clinkers)

KITSIS, S.I., inzh.

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(Hoisting machinery) (Automatic control)

KITSIS, S.I., inzh.

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short-term mode of operation. Elektrotehnika 35 no.1:63
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KITSIS, S.I., inah.

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energ. 18 no.8:20-21 Ag '63. (MIRA 16:9)
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22 no.3:13-14 Mr '63. (MIRA 16:4)
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KITSIS, S.I., inzh.

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17 no.12:47-48 D '62. (MIRA 17:4)

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KITSKIY, B.

USSR/Gasoline
Engines, Gasoline

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New Model Cars," B. Kitskiy, Engr, 2 pp

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<p>ASS. 554 METALLOGICAL LITERATURE CLASSIFICATION</p>	
<p>1000 1000000</p>	<p>1000 1000000</p>

VOINOV, N. P.; KONEV, B. P.; KITSKIY, B. P.
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Toplivo i Smazka Otechestvennykh Legkovykh Avtomobilei (Fuel and Oil for
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Toplivo i smazka otchestvennykh bol'shegruznykh avtomobiley (Fueling and lubrication of Russian heavy trucks, by) B. P. Kitskiy (and) M. S. Khanin, Moskva, Gostoptekhizdat, 1952.

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SO: 3/5

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Effect of the temperature of combustion chamber walls on operating
conditions of the IAAS-204 engine. [Trudy] NVTU no. 83:172-181 '58.
(Gas and oil engines) (MIRA 11:6)

ZAYTSEV, V.; LERNER, M.; GUREYEV, A.; ZABRYANSKIY, Ye.; KITSKIY, B.

New antiknock compound for gasolines. Avt.transp. 40 no.1:17-19
Ja '62. (MIRA 15:1)
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ANDRONIKASHVILI, E.L., akademik; BUDA, B.G.; DEVNOZASHVILI, D.S.;
KIKHADZE, G.I.; KITSHARISHVILI, E.S.; TOPSHYAN, L.S.;
CHANTURIYA, V.M.

Low-temperature loop of an IRT-2000 reactor. Soob. AN Grus.
SSR 34 no.1:45-52 Ap'64 (MIRA 17:7)

1. AN Gruzinskoy SSR (for Andronikashvili).

KITSNIK, A., KOCH, R., kand.tekhn.nauk

Organic matter and mineral contents in the various granulometric
classes of pulverized kukersite. Izv. AN Est. SSR, Ser. fiz.-mat.
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1. Institut khimii A' Estonskoy SSR.

KITSOVSKAYA, I.A.

Study of correlation between the basic neural processes in
rats under the effect of SHF of various intensity. Trudy
Inst. gig. truda i prof. AMN SSSR no.1:75-80 '60.

(MIRA 16:12)

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GORDON, Z.V., KITOVSKAYA, I.A., and TOLOSKAYA, M.S.

All at the Institute of Labor Hygiene and Professional Diseases,
Academy of Medical Sciences USSR, Moscow - "Biological effects of
microwaves of small intensity"

Report to be submitted for the 4th Intl. Conf. on
Medical Electronics, New York, N.Y., 16-21 July 1961

1. 35862-66 EWT(1) DO	
ACC NR: AP6022515 (V)	SOURCE CODE: UR/0391/66/000/007/0005/0009 SD
AUTHOR: <u>Fukalova, P. P. (Moscow); Tolgskaya, M. S. (Moscow);</u> <u>Nikogosyan, S. V. (Moscow); Kitsovskaya, I. A. (Moscow); Zenina, I. N. (Moscow)</u>	
ORG: <u>Institute of Industrial Hygiene and Occupational Diseases, AMN</u> <u>SSSR (Institut gigiyeny truda i profzabolevaniy AMN SSSR)</u>	
TITLE: <u>Research data on the standardisation of EMF's in the short- and</u> <u>ultrashort-wave ranges</u>	
SOURCE: <u>Gigiyena truda i professional'nyye zabolevaniya, no. 7,</u> <u>1966, 5-9</u>	
TOPIC TAGS: <u>microwave, microwave biologic effect, central nervous</u> <u>system, UHF, human physiology, animal physiology, animal experiment,</u> <u>industrial hygiene</u>	
ABSTRACT: <u>In a survey of radio and television stations and establish-</u> <u>ments which process dielectrics thermally, measurements using</u> <u>various dosimeters showed that field intensities around short- and</u> <u>ultrashort-wave sources were 8-450 v/m and 4-220 v/m, respectively.</u> <u>The reaction speed and accuracy were studied in personnel exposed to</u>	
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ACC NR: AP6022515

these sources. Chronometric tests, observations of visual motor reactions, and work capacity were studied and are summarized in Table 1.

Table 1. Changes in the physiological functions of duty personnel in two radio stations (%)

Observed parameter	Shift									
	Evening		Day		Night					
					Station I			Station II		
	Station I	Station II	Station I	Station II	Before sleep	After sleep	End of shift	No sleep		
Latent period of positive conditioned reflexes	0.0	0.0	16.1	12.6	2.2	2.2	6.6	0.0	11.1	42.5
Subsequent inhibition	6.9	4.1	24.3	18.7	0.0	2.1	13.0	0.0	8.1	17.3
Disinhibition of differentiation	20.0	10.0	40.0	20.0	60.0	60.0	70.0	30.0	2.1	60.0
Steadiness of attention	21.4	33.0	60.0	21.4	33.0	19.9	23.3	25.0	50.0	36.6
No. of errors	1.5	1.0	2.5	1.0	3.0	3.0	4.0	2.0	2.5	4.0

It was concluded that the action of EMP's is aggravated by inefficient work-rest cycles which lead to shifts in various physiological reactions

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ACC NR: AP6022515

and the development of fatigue. To study the mechanism of the effect of EMF's, animals were exposed to 14, 88, and 69.7 Mc fields (5000 v/m intensity). Animals in the ultrashort-wave range were killed within 5 min, while those in the shortwave range died in 1 hr and 40 min. Nonthermal (no integral thermal effect such as increased body temperature) threshold intensity for the ultrashort-wave range was 150 v/m and for the shortwave range, 2250 v/m. Chronic exposure to these intensities (plus exposure to an ultrashort-wave intensity of 300 v/m) did not result in any substantial changes in body weight dynamics compared to control animals. However, a decrease in brain-stem cholinesterase activity occurred more rapidly during exposure to ultrashort waves than exposure to short waves. Both regimens decreased the excitability and weakened the inhibition process in chronically exposed rats. Such exposure also tended to depress brain biopotentials progressively. An ultrashort-wave intensity of 10 v/m and shortwave intensity of less than 50 v/m is a subthreshold irritant. Photos show the results of a cytomorphological examination of neural structures in exposed animals which revealed thickening of neural fibers, swelling and vacuolization of cell protoplasm in the thalamo-hypothalamic area and medulla oblongata, and local karyocytolysis of individual neurons. Shriveling of individual pyramid cells and individual vacuoles in neurons of the brain cortex was noted. Thus, it was found that an ultrashort-wave intensity of 150 v/m and a shortwave intensity of 2250 v/m is more than

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L 35862-66

ACC NR: AP6022515

sufficient to cause physiological and morphological changes in neural structures and that the intensity of the effect depends on wave range, field intensity, and exposure duration. On the basis of these data, a permissible exposure intensity of 5 v/m is recommended for workers around ultrashort-wave sources and 20 v/m for workers around shortwave sources. Orig. art. has: 6 figures and 1 table. [CD]

SUB CODE: 06/ SUBM DATE: 07Jan65/ ORIG REF: 017/ OTH REF: 001
ATD PRESS: 50 37

Card 4/4 *lll*

L 47127-66 EWT(d) IJP(o)

ACC NR: AR6016019

SOURCE CODE: UR/0271/66/000/001/B003/B003

AUTHOR: Kisel', V. A.; Kitsul, I. V.

50
B

TITLE: Trigonometric interpolation with multiple equidistant nodal points

SOURCE: Ref. zh. Avtomat. telemekh. i vychisl. tekhn., Abs. 1B20

REF SOURCE: Tr. uchebn. in-tov svyazi. M-vo svyazi SSSR, vyp. 24, 1965, 169-176

TOPIC TAGS: polynomial, interpolation, harmonic function, approximation method

ABSTRACT: A method is analyzed for plotting the trigonometric polynomials which for interpolating the values of a given function and values of its derivatives in equidistant modal points. Three cases are analyzed: when an interpolating function is of even parity, odd parity, and random. The given solution can be used for calculating the harmonic corrections, vocoders, and synthesizers, as well as equipment, the calculation of which is based on an approximation of the given function and its derivatives of trigonometric polynomials. Orig. art. has: 1 figure.

[Translation of abstract]

[NT]

SUB CODE: 12/

Cord 1/1 afa

UDC: 681.142.33.001

ALDEN, V.I.

Geochronological significance of the ultrabasic alkali intru-
sions in the Alden field. Trudy lab. geol. dokw. no.19:
228-235 '64 (MIRA 1718)

KITSUL, V.I.

Metamorphism of carbonate rocks of the Ladoga formation. Trudy lab.
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BOZHKOV, I.S., KITSUL, V.I.

**A deposit of platinum in the Aldan shield. Geol. rad. masterosh.
no. 4,74-84 JI-Ag '60. (MIRA 13:8)**

- 1. Yakutskiy filial Sibirskogo otdeleniya AN SSSR.
(Aldan Plateau--Platinum)**

ROZHKOV, I.S., glav. red.; KITSUL, V.I., kand. geol.-min. nauk,
otv. red.

[Petrography of the metamorphic and igneous rocks of the
Alden Shield] Petrografiia metamorficheskikh i izverzhennyykh
porod Aldanskogo shchita. Moskva, Nauka, 1964. 163 p.
(MIRA 17:6)

1. Akademiya nauk SSSR. Yakutskiy filial, Yakutsk. Institut
geologii. 2. Chlen-korrespondent AN SSSR (for Rozhkov).

KITSUL, V.I.

Grades of the progressive regional metamorphism of carbonate rocks
in the Ladoga formation. Trudy Lab. geol. dokl. no.11:230-239
160. (MIRA 14:1)

(Ladoga Lake region--Rocks, Carbonate)

KITSUL, V.I.; BOGOMOLOV, M.A.

Concerning G.V.Andreev's article "Contact-infiltration skarns near
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26 no.1:99-100 Ja '61. (MIRA 1516)
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KALANTAROV, A.P., red. izd-va; GUSEVA, A.P., tekhn. red.

[Petrology of carbonate rocks in the Ladoga formation] Petrologiia karbonatnykh porod Ladozhskoi formatsii. Moskva, Izd-vo Akad. nauk SSSR, 1963. 170 p. (MIRA 16:5)
(Ladoga Lake region--Rocks, Carbonate)

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Experimental data on a new signification of the activity of
cholinesterases, Studii cerc fisiol 6 no.1:19-26 '61,
(REAI 10:9)

1. Institutul oncologic, Bucuresti.

(ENZYMES) (CHOLINESTERASES)